

Lecture 20 - 11/17

Announce:

- Will do group theory/symm. for rest of class \rightarrow who knows?
I'll post pages from Dresselhaus book.

Program:

- Many-electron atoms
- Symmetry

c) Hund's rules

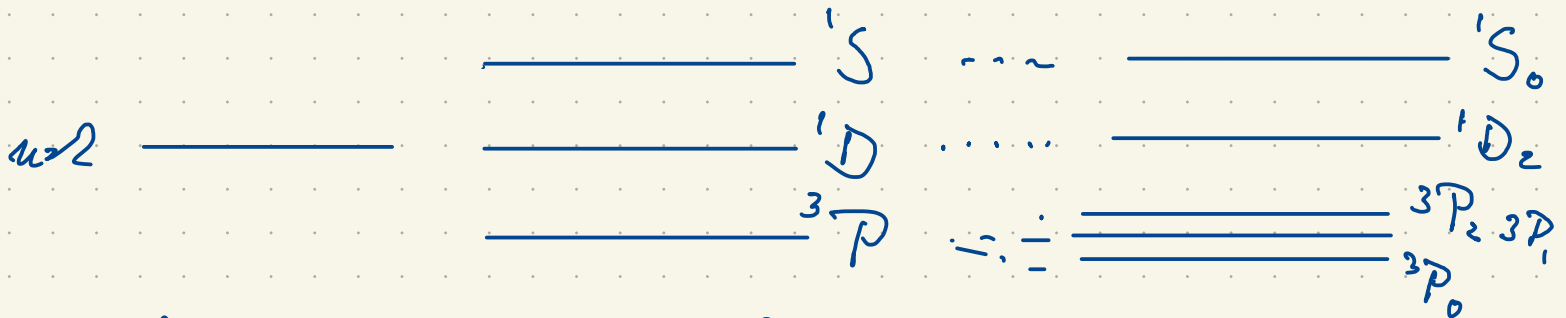
① In any given subshell (same n , same L), state with largest S is energetically lowest.

② For same shell (same n), same S state with highest L has lowest energy.

Reason in both cases: get e's as far apart as possible!

① gives bigger energy difference than ②.

Example: $1s^2 2s^2 2p^2$



Why? Open sub-shell: $2p^2$

$$|l_1 - l_2| \leq L \leq |l_1 + l_2| \quad (l_1 = l_2 = 1)$$

$$0 \leq L \leq 2$$

$$|s_1 - s_2| \leq S \leq |s_1 + s_2| \quad (s_1 = s_2 = \frac{1}{2})$$

$$0 \leq S \leq 1$$

\Rightarrow in principle: $L = \{0, 1, 2\}$, $S = \{0, 1\}$

${}^{2S+1}L_J \rightarrow {}^1S_0, {}^3S_1, {}^1P_1, {}^3P_0, {}^1D_2, {}^3D_3$

which allowed?

Symmetry: S : $S=0$ (odd) $S=1$ (even)

L : S, D (even) P (odd)

only \Rightarrow ${}^1S_0, {}^1D_2, {}^3P_0$ allowed!

1st Hund's rule: 3P lowest

2nd " : S higher than D

How to sort different J:

in principle: case-to-case

in practice:

"regular" (Energy $\sim J$) - for up to half-filled subshells (✓)

"inverted" ordering: otherwise

Note: "reason" for Hund's rules:

①: subshells (e.g. p_x, p_y, p_z) singly occupied before pairing (due to e^- -repulsion), all the same spin (if possible) \Leftrightarrow more effective screening

②: higher L \Leftrightarrow large distance between e^-

③: $\langle H_{\text{spin-orbit}} \rangle = \zeta \langle \vec{L} \cdot \vec{S} \rangle \Rightarrow E_{J+1} - E_J \Big|_{\substack{\text{equal} \\ L, S}} = \zeta (J+1)$
 $\zeta > 0$ ("regular") for subshell ruled by e^-
 $\zeta < 0$ ("inverted") holes

in general: (p^1 and p^5) or (p^2 and p^4)
(d^1 and d^9) or (d^2 and d^8)...

have same configuration (i.e. e^- in empty subshells have same effect as holes in full subshell!)

9) Symmetry

a) Short review (selection rules)

So far:

parity: $\Delta l, \Delta j = \text{odd}$

spherical: $\Delta l = \pm 1, \Delta j = \pm 1, 0$

symmetry: $\Delta m_l, \Delta m_j = \pm 1, 0$

In general:

System is symmetric under operator \hat{T} :



$$[H, \hat{T}] = 0$$

b) Time reversal symmetry

$$\hat{T}: t \rightarrow -t$$

even: $\vec{r} \Rightarrow \vec{r}, \vec{F}, E, V(\propto r^k), \vec{d}, \dots$

odd: $\dot{\vec{r}} \Rightarrow \vec{L}, \vec{f}, \vec{B}, \dots$

Simple case: (S.Eq.)

$$\hat{T} \psi(t) = \hat{T} \psi(0) e^{-\frac{i}{\hbar} H t} = \psi(0) e^{+\frac{i}{\hbar} H t} = \psi^*(t)$$

Properties

$$\textcircled{1} [\hat{T}, H] = 0$$

$$\textcircled{2} \hat{T} i = -i \hat{T} \Rightarrow \hat{T} a = a^* \hat{T}$$

"anti-linear"

$$\textcircled{3} \hat{T} \psi(t) = \psi^*(t) \hat{T}$$

④ no spin: \hat{T} is just "c.c." operator \hat{K}

with spin: $\hat{T} = \hat{K} \sigma_y$

no spin: $\hat{T}^2 = 1$

with spin: $\hat{T}^2 = -1$

$\vec{d} \cdot \vec{B}$

(more details: Dresselhaus script)

Consequences:

- \hat{T} (and \hat{K}) is "anti-unitary"
- Kramer's theorem: for $\hat{T}^2 = -1$ has to have at least a 2D ^{irred} representation (two degenerate eigenfunctions)

(Proof: time-reversed state also eigenstate, but cannot be the same.)

- Electric field cannot split states $\pm |m\rangle$: \vec{E} is even under \hat{T} , angular mom. is odd: energy has to be invariant under \hat{T} . □

- Electric dipole moment: EDM:

An intrinsic (i.e. not induced) EDM violates both T and P symmetries.

Proof: total angular mom (of particle, e.g. e^-): \vec{J}
assume dipole moment \vec{d} :

\vec{d} has to be parallel to \vec{J} (all components \perp to \vec{J} average out, \vec{J} only preferred direction!)
 $\Rightarrow \vec{d} = \beta \vec{J} \quad (\beta \in \mathbb{C})$

\vec{J} changes under \hat{T} , but not under \hat{P} .

\vec{d} changes under \hat{P} , but not under \hat{T}

\Rightarrow either (separably):

$$\vec{d} = \beta \vec{J} \quad \stackrel{\hat{T}}{\Rightarrow} \quad \vec{d} = -\beta \vec{J}$$

$$\Rightarrow \beta = 0 \quad \Rightarrow \quad \vec{d} = \vec{0} \quad \square$$