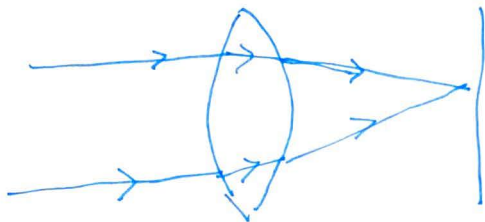
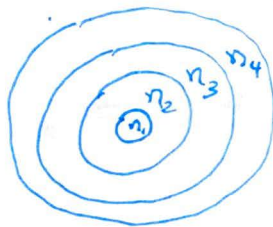


Geometrical Optics

Light before Maxwell - Rays.



Also an actual optical fiber is graded.



$$n_1 > n_2 > n_3 \dots$$
$$n(r)$$

We assume $n(r), \epsilon(r)$ vary on a scale
 $\downarrow \Rightarrow \lambda$, wavelength of light.

Analysis of WKB in quantum mechanics

$$\vec{E}(\vec{r}, t) = \vec{E}_0 e^{i\omega S(r)/c} e^{-i\omega t}$$

$$\vec{B}(\vec{r}, t) = \vec{B}_0 e^{i\omega S(r)/c} e^{-i\omega t}$$

$S(\vec{r})$ is the "eikonal"

\Leftrightarrow WKB in QM.

Then Maxwell's equations become

$$\boxed{\vec{E}_0 \cdot \vec{\nabla} S = 0} \quad (\nabla \cdot D = 0)$$

We ignore $\vec{\nabla} \epsilon$ because $d \rightarrow \lambda$

$$\boxed{\vec{B}_0 \cdot \vec{\nabla} S = 0} \quad (\nabla \cdot B = 0)$$

$$(\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t})$$

$$\frac{1}{c} \vec{\nabla} S \times \vec{E}_0 = \vec{B}_0$$

$$\frac{1}{c} \vec{\nabla} S \times \vec{B}_0 = -\mu \epsilon \vec{E}_0 \quad (\vec{\nabla} \times \vec{B} = \mu \epsilon \frac{\partial \vec{E}}{\partial t})$$

Consistency of last 2 equations leads to

$$\vec{\nabla} S \cdot \vec{\nabla} S = c^2 \mu \epsilon(\mathbf{r})$$

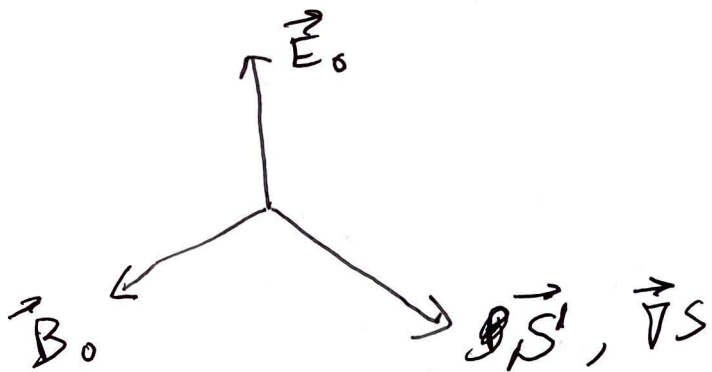
$$\boxed{\vec{\nabla} S \cdot \vec{\nabla} S = n^2(\mathbf{r})}$$

Eikonal equation.

$n(\mathbf{r}) \rightarrow$ refractive index.

Poynting vector

$$\vec{S} = \frac{1}{\mu} \vec{E} \times \vec{B} = \frac{1}{2\mu c} |E_0|^2 \vec{\nabla} S$$

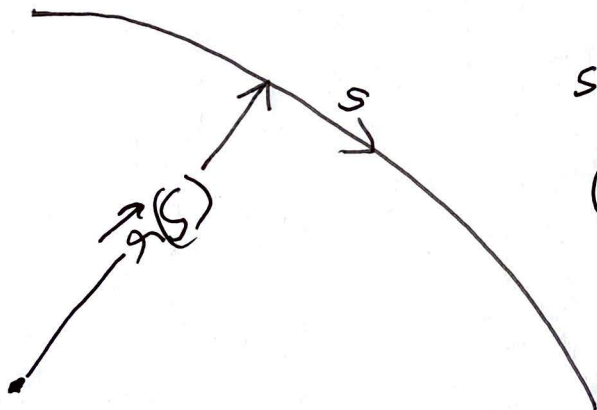


We write

$$\vec{v}_S = n(\vec{r}) \hat{k}(\vec{r})$$

$\hat{k}(\vec{r}) \rightarrow$ unit vector in the direction of the ray

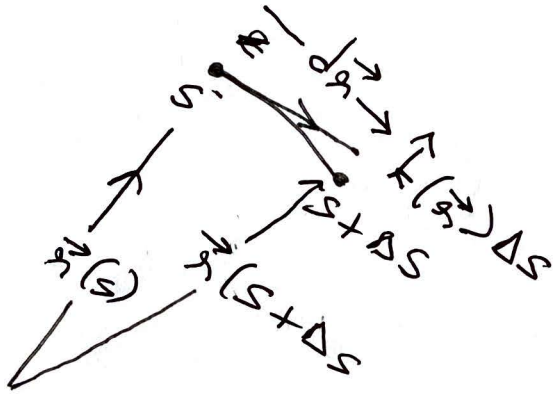
Path of ray



$s \rightarrow$ path length along ray
(analogous to proper time in relativity)

$$\text{So } d\vec{r}^2 = ds^2$$

i.e. $\left| \frac{d\vec{r}}{ds} \right| = 1.$



We have $d\vec{r} = \hat{k}(\vec{r}) \Delta S$

Or $\frac{d\vec{r}}{dS} = \hat{k}(\vec{r})$ (OK because $\left| \frac{d\vec{r}}{dS} \right| = |\hat{k}(\vec{r})| = 1$)

So $\vec{\nabla} S = n(\vec{r}) \frac{d\vec{r}}{dS}$

Therefore $\frac{d}{dS} \left(n(\vec{r}) \frac{d\vec{r}}{dS} \right) = \vec{\nabla} \frac{dS}{dS}$

But $\frac{d}{dS} = \hat{k} \cdot \vec{\nabla}$ So $\frac{dS}{dS} = \hat{k} \cdot \vec{\nabla} S$
 $= \hat{k} \cdot \hat{k} n(\vec{r})$

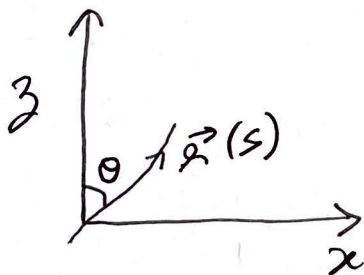
So finally we obtain the ray equation

$$\frac{d}{ds} \left[n(\vec{r}) \frac{d\vec{r}}{ds} \right] = \vec{\nabla} n(\vec{r})$$

(i) Allows solutions with $\left| \frac{d\vec{r}}{ds} \right| = 1$ for all s .

(ii) Generalization of Snell's law.

Suppose $\vec{\nabla} n(\vec{r})$ is a vector in the \hat{z} direction.



Then $n(\vec{r}) \frac{dx}{ds} = \text{constant}$

$$\text{or } \boxed{n \sin \theta = \text{constant}}$$

Why is $\left| \frac{d\vec{r}}{ds} \right| = 1$ for ray equation?

$$\frac{d}{ds} \left(\frac{d\vec{r}}{ds} \right)^2 = 0$$

$$\Rightarrow \frac{d^2 \vec{r}}{ds^2} \cdot \frac{d\vec{r}}{ds} = 0$$

The ray equation preserves this condition

LHS =

$$\frac{d}{ds} \left[n(\vec{r}(s)) \frac{d\vec{r}}{ds} \right] = \left(\vec{\nabla} n \cdot \frac{d\vec{r}}{ds} \right) \frac{d\vec{r}}{ds} + n(\vec{r}) \frac{d^2 \vec{r}}{ds^2}$$

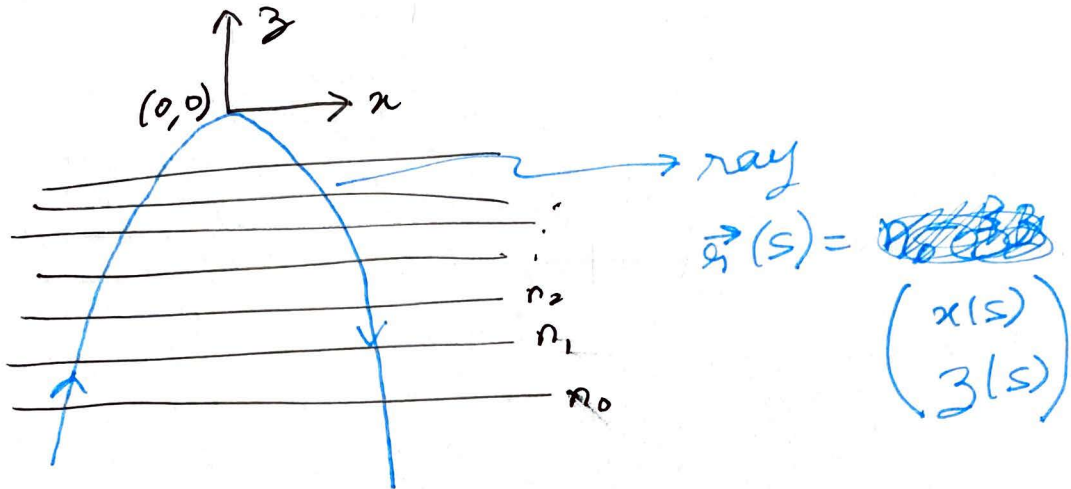
RHS = $\vec{\nabla} n(\vec{r}) \cdot \frac{d\vec{r}}{ds}$

$$\begin{aligned} (\text{LHS} - \text{RHS}) \cdot \frac{d\vec{r}}{ds} &= \left(\vec{\nabla} n \cdot \frac{d\vec{r}}{ds} \right) \frac{d\vec{r}}{ds} \cdot \frac{d\vec{r}}{ds} + n(\vec{r}) \frac{d^2 \vec{r}}{ds^2} \cdot \frac{d\vec{r}}{ds} \\ &\quad - \frac{d\vec{r}}{ds} \cdot \vec{\nabla} n \\ &= n(\vec{r}) \frac{d^2 \vec{r}}{ds^2} \cdot \frac{d\vec{r}}{ds} = 0 \end{aligned}$$

QED

Reflection by a graded dielectric

$$n(\vec{r}) = n_0 - \beta z$$



Ray Equation

$$\frac{d}{ds} \left[\vec{n}(\vec{r}(s)) \frac{d\vec{r}}{ds} \right] = -\beta \begin{pmatrix} 0 \\ 1 \end{pmatrix} \quad (*)$$

Boundary condition $\vec{r}(0) = \begin{pmatrix} 0 \\ 0 \end{pmatrix}$

and $\frac{d\vec{r}}{ds}(0) = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$.

Integrating (*) and imposing boundary conditions

$$\frac{dx}{ds} = \frac{n_0}{n_0 - \beta z} \quad ; \quad \frac{dz}{ds} = \frac{-\beta s}{n_0 - \beta z}$$

Solution for $z(s)$

$$z(s) = \frac{1}{\beta} \left[n_0 - \sqrt{n_0^2 + \beta^2 s^2} \right]$$

Plug into equation for $x(s)$ and solve

$$x(s) = \frac{n_0}{\beta} \ln \left(\frac{\beta s + \sqrt{n_0^2 + \beta^2 s^2}}{n_0} \right)$$

Verify that (magically?)

$$\left(\frac{dx}{ds} \right)^2 + \left(\frac{dz}{ds} \right)^2 = 1.$$